## Chapter 8

## Representations of two compact groups.

### 8.1 Representations of the permutation group

Permutations on $n$ symbols is the set of one to one mappings $\sigma$ of a set $X=\left\{x_{1}, x_{2}, \ldots, x_{n}\right\}$ of $n$ elements on to itself. It is a finite group $G$ with $n!$ elements. Given $\sigma \in G$ we can look at the orbits of $\sigma^{n} x$ and it will partition the speace $X$ into orbits $A_{1}, A_{2}, \ldots, A_{k}$ consisting of $n_{1} \geq \ldots \geq n_{k}$ points so that $n=n_{1}+\cdots+n_{k}$. If $\hat{\sigma}=s \sigma s^{-1}$ is conjugate to $\sigma$ then the orbits of $\hat{\sigma}$ will be $s A_{1}, s A_{2}, \ldots, s A_{k}$ so that the partition $n_{1} \geq \ldots \geq n_{k}$ of $n$ into $k$ numbers will be the same for $\sigma$ and $\hat{\sigma}$. Conversely if $\sigma$ and $\hat{\sigma}$ have the same partition of $n$, then the orbits $A_{1}, \ldots, A_{k}$ and $B_{1}, \ldots, B_{k}$ can be so arranged that the corresponding orbits have the same cardinality. We can find $s \in G$ that maps $A_{i} \rightarrow B_{i}$ in a one-to-one and onto manner and we can reduce the problem to the case where $A_{i}=B_{i}$ for every $i$. We can now relablel the points with in each $A_{i}$ i.e. find a permutation of $n_{i}$ elements, such that both $\sigma$ and $\hat{\sigma}$ look the same on each $A_{i}$. We can now prove the following theorem.

Theorem 8.1. The number of distinct inequivalent irreducible representations of $G$ is the same as the number of distinct partitions $\mathcal{P}(n)$ of the integer $n$.

Proof. The space $\mathcal{U}$ of functions $u$ that are invariant under conjugation i.e satisfy $u\left(h^{-1} g h\right)=u(g)$ is spanned by characters and its dimension equals the number of conjugacy classes, i.e the number of distinct sequences $n_{1} \geq$
$n_{2} \geq \cdots \geq n_{k}>0$ such that $n=n_{1}+\cdots+n_{k}$. This is the same as the number of distinct partitions of $n$.

We will now construct a distinct irreducible representation for each partition of $n$. Given a partition $\lambda$ of $n$ into $n_{1} \geq n_{2} \geq \cdots \geq n_{k}$, we associate a diagram called the Young diagram corresponding to $\lambda$. It looks like

when $n=9$ and the partition is $3,3,2,1$. Each set in the partition is a row consisting of squares whose number is the number of elements in that set. The rows are arrnged in decreasing order of their size and aligned on the left. A Young tableau $\mathbf{t}$ is a diagram $\lambda$ with the boxes filled in arbitrarily by the numbers $1,2, \cdots, 9$, like

| 4 | 1 | 8 |
| :--- | :--- | :--- |
| 5 | 2 | 7 |
| 3 | 6 |  |
| 9 |  |  |
|  |  |  |
|  |  |  |

The rows of the array are of length $n_{1} \geq \cdots \geq n_{k}$. For any diagram there are $n$ ! tableaux. The tableaux $\mathbf{t}$ above has rows $[4,1,8],[5,2,7][3,6]$ and [9]. A tabloid is when the order of entries with in a row is not relevant. The tabloid $\tau=\{\mathbf{t}\}$ consists of $\{1,4,8\},\{2,5,7\},\{3,6\},\{9\}$. There are $\frac{n!}{n_{1}!\cdots n_{k}!}$ tabloids corresponding to a diagram $\lambda$ corresponding to the partition $n_{1} \geq$ $n_{2} \geq \cdots \geq n_{k}$. The subgroup $C_{\mathbf{t}}$ of the permutation group consists of permutations within each column. In our case it consists of $4!\times 3!\times 2!=288$ elements. An arbitrary permutation of $4,5,3,9$, an arbitrary permutation of $1,2,6$ and one of 8,7 . For any permutation $s, \sigma(s)= \pm 1$ is the sign of the permutation. We define an abstract inner product space $V-\lambda$ of dimension $\frac{n!}{n_{1}!\cdots n_{k}!}$ with the orthonormal basis $e_{\tau}$ as $\tau$ varies over all the tabloids of the diagram $\lambda$. We define $n$ ! vectors $f_{\mathbf{t}}$ in $V_{\lambda}$ as $\mathbf{t}$ varies over the set $\mathbf{T}$ of tableau of the diagram $\lambda .\{\mathbf{t}\}$ is the tabloid that is the equivalence class of all the tableaux that are obtained by permutations with in rows of the tableaux t. The group element, namely a permutation $s$ acts naturally on the set of
tableaux $\mathbf{T}$. If two tableaux $\mathbf{t}_{1}$ and $\mathbf{t}_{2}$ belong to the same tabloid so do $s \mathbf{t}_{1}$ and $s \mathbf{t}_{2}$. Therefore $s$ acts on the set $\mathcal{T}$ of tabloids as well. Define

$$
f_{\mathbf{t}}=\sum_{s \in C_{\mathbf{t}}} \sigma(s) e_{s\{\mathbf{t}\}}
$$

The $\left\{f_{\mathbf{t}}: \mathbf{t} \in \mathbf{T}\right\}$ may not be linearly independent and the span of $\left\{f_{\mathbf{t}}\right\}$ is denoted by $W$. One defines a representation of $\pi_{\lambda}(g)$ of the permutation group on $V_{\lambda}$ and $W_{\lambda}$ corresponding to the diagram $\lambda$ by defining

$$
\pi_{\lambda}(g) f_{\mathbf{t}}=g f_{\mathbf{t}}=\sum_{s \in C_{\mathbf{t}}} \sigma(s) e_{g s\{\mathbf{t}\}}
$$

Theorem 8.2. Each $\pi_{\lambda}$ is irreducible. For two distinct diagrams they are inequivalent. We therefore have all the representations.

Proof is broken up into lemmas.

## Lemma 8.3.

$$
\pi_{\lambda}(g) f_{\mathbf{t}}=f_{g \mathbf{t}}
$$

Proof.

$$
\begin{aligned}
\sum_{s \in C_{\mathbf{t}}} \sigma(s) e_{g s\{\mathbf{t}\}} & =\sum_{s \in C_{\mathbf{t}}} \sigma(s) e_{g s g^{-1} g\{\mathbf{t}\}} \\
& =\sum_{g s g^{-1} \in C_{g \mathbf{t}}} \sigma(s) e_{g s g^{-1} g\{\mathbf{t}\}} \\
& =\sum_{s^{\prime} \in C_{g \mathbf{t}}} \sigma\left(s^{\prime}\right) e_{s^{\prime} g \mathbf{t}}=f_{g \mathbf{t}}
\end{aligned}
$$

Lemma 8.4. Suppose $\lambda, \mu$ are two different diagrams, $\mathbf{t} a \lambda$-tableaux and $\mathbf{t}^{*}$ a $\mu$-tableaux. Suppose

$$
\sum_{s \in C_{\mathbf{t}}} \sigma(s) e_{s\left\{\mathbf{t}^{*}\right\}} \neq 0
$$

Then $n_{1} \geq m_{1}, n_{1}+n_{2} \geq m_{1}+m_{2}, \cdots$ where $n_{1}, \ldots, n_{k}$ and $m_{1}, \ldots, m_{\ell}$ are the two partitions corresponding to $\lambda$ and $\mu$. We say then that $\lambda \geq_{1} \mu$. If $\lambda=\mu$ then the sum is $\pm f_{\mathrm{t}}$.

Proof. Suppose that two elements $a, b$ are in the same row of $\mathbf{t}^{*}$ and in the same column of $\mathbf{t}$. Then the permutation $p=\{a \leftrightarrow b\}$ is in $C_{\mathbf{t}}, e_{s p\{\mathbf{t}\}}=e_{s\{\mathbf{t}\}}$, with $\sigma(s p)+\sigma(s)=0$. So the sum adds up to 0 which is ruled out. Hence, no two elements in the same row of $\mathbf{t}^{*}$ can be in the same column of $\mathbf{t}$. In particular $\mathbf{t}$ must have at least as many columns as the number of elements in the first row of $\mathbf{t}^{*}$ proving $n_{1} \geq m_{1}$. A variant of this argument proves $\lambda \geq_{1} \mu$. Suppose now that $\lambda=\mu$. Again all the elements in any row of $\mathbf{t}^{*}$ appear in different columns of $\mathbf{t}$. Then there is a permutation $s^{*} \in C_{\mathbf{t}}$ such that $s^{*} \mathbf{t}=\mathbf{t}^{*}$. The sum is unaltered if we replace $\mathbf{t}$ by $s^{*} \mathbf{t}$ except for the factor of $\sigma\left(s^{*}\right)= \pm 1$.

Lemma 8.5. Let $u \in V$ corresponding to a diagram $\lambda$. Let $\mathbf{t}$ be any $\lambda$ tableau. Then

$$
\sum_{s \in C_{\mathbf{t}}} \sigma(s) \pi_{\lambda}(s) u=c f_{\mathbf{t}}
$$

Proof. $u$ is a linear combination of $e_{\left\{\mathbf{t}^{*}\right\}}$ for different tabloids $\left\{\mathbf{t}^{*}\right\}$ corresponding to $\lambda$. Each one from the previous lemma yields $c\left(\mathbf{t}^{*}\right) f_{\mathbf{t}}$ with $c\left(\mathbf{t}^{*}\right)=0, \pm 1$. Add them up!

Let us define

$$
A_{\mathbf{t}}=\sum_{s \in C_{t}} \sigma(s) \pi_{\lambda}(s)
$$

acting on $V$. We already have an inner product that makes $\pi(s)$ orthogonal.

$$
\begin{aligned}
<A_{\mathbf{t}} u, v> & =\sum_{s \in C_{\mathbf{t}}} \sigma(s)<\pi(s) u, v> \\
& =\sum_{s \in C_{\mathbf{t}}} \sigma(s)<u, \pi\left(s^{-1}\right) v> \\
& =\sum_{s \in C_{\mathbf{t}}} \sigma(s)<u, \pi(s) v> \\
& =<u, A_{\mathbf{t}} v>
\end{aligned}
$$

because $\sigma\left(s^{-1}\right)=\sigma(s)$.
Lemma 8.6. If $U_{\lambda}$ is any invariant subspace of $V$ then either $U_{\lambda} \supset W_{\lambda}$ or $U_{\lambda} \perp W_{\lambda}$. This proves the irreducibility of $W_{\lambda}$.

Proof. Suppose for some $u \in U_{\lambda} \subset V_{\lambda}$ and $\mathbf{t}$ is a $\lambda$-tableau. We saw that $A_{\mathbf{t}} u=c_{\mathbf{t}} f_{\mathbf{t}}$ for some constant $c_{\mathbf{t}}$. Suppose for some $u \in U_{\lambda}$ and $\mathbf{t}, c_{\mathbf{t}} \neq 0$. Then $f_{\mathbf{t}}=\frac{1}{c_{\mathrm{t}}} A_{\mathbf{t}} u \in U_{\lambda}$. Since $\pi(g) f_{\mathbf{t}}=f_{g \mathbf{t}}$ and $\left\{f_{g \mathbf{t}}\right\}$ span $W_{\lambda}$, it follows that $W_{\lambda} \subset U_{\lambda}$. If $c_{\mathbf{t}}=0$ for all $u$ and $\mathbf{t}, A_{\mathbf{t}} u=0$ and it follows that

$$
<u, f_{\mathbf{t}}>= \pm<u, A_{\mathbf{t}} f_{\mathbf{t}}>= \pm<A_{\mathbf{t}} u, f_{\mathbf{t}}>=0
$$

and $u \in W^{\perp}$.
Lemma 8.7. Let $T$ intertwine the representations on $V_{\lambda}$ and $V_{\mu}$. Suppose $W_{\lambda}$ is not contained in Ker $T$. Then $\lambda \geq_{1} \mu$.

Proof. Ker $T$ is invariant under $\pi(g)$ and if it does not contain $W \lambda$ it is orthogonal to it.

$$
0 \neq T f_{\mathbf{t}}=T A_{\mathbf{t}} f_{\mathbf{t}}=A_{\mathbf{t}} T f_{\mathbf{t}}
$$

$T f_{\mathbf{t}}$ is a combination of $e_{\{\mathbf{t}\}}$ of $\mu$-tableau $\{\mathbf{t}\}$. So at least one of them $A_{\mathbf{t}} e_{\{\mathbf{t}\}}$ is nonzero forcing $\lambda \geq_{1} \mu$.

Lemma 8.8. If $T \neq 0$ intertwines $W_{\lambda}$ and $W_{\mu}$, then $\lambda=\mu$.
Proof. Extend $T$ by making it 0 on $\left(W_{\lambda}\right)^{\perp}$ and we see that $\lambda \geq_{1} \mu$. The argument is symmetric.

### 8.2 Representations $\mathrm{SO}(3)$

We will consider the irreducible representations of the group $G$ of rotations in $R^{3}$. These are orthogonal transformations of determinant 1, i.e. that preserve orientation. An element $g \in G$ is represented as the matrix

$$
\left[\begin{array}{ccc}
t_{1,1}(g) & t_{1,2}(g) & t_{1,3}(g) \\
t_{2,1(g)} & t_{2,2}(g) & t_{2,3}(g) \\
t_{3,1(g)} & t_{3,2}(g) & t_{3,3}(g)
\end{array}\right]
$$

There is the trivial representation $\pi_{0}(g) \equiv I$. Then there is a natural three dimensional representation where $\pi_{1}(g)=t(g)=\left\{t_{i, j}(g)\right\}$ and it can be viewed as a unitary representation in $\mathcal{C}^{3}$. This representation is irreducible and faithful, i.e. it separates points of $G$.

As we saw in the general theory, the characters can be used to identify the irreducible representations. It helps to know what the conjugacy classes are. Given two orthogonal matrices $g_{1}$ and $g_{2}$, when can we find a $g$ such that $g g_{1} g^{-1}=g_{2}$ ? The eigen values of $g_{1}$ are $1, e^{ \pm i \theta_{1}}$ and therefore in order for $g_{1}$ and $g_{2}$ to be mutually conjugate we neeed $\theta_{1}= \pm \theta_{2}$ or $\cos \theta_{1}=\cos \theta_{2}$. Conversely one can show that that if $g_{1}$ and $g_{2}$ have the same eigenvalues then they are indeed conjugate. If we use a $g$ to align the eigenspace correponding to 1 , then we need to show essentially that rotation by $\theta$ and $-\theta$ are conjugate. We can use the matrix

$$
\left[\begin{array}{ccc}
-1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & -1
\end{array}\right]
$$

to achieve this.
We will use the infinitesimal method to study irreducible representations. If $A=\left\{a_{i, j}\right\}$ is a real skewsymmetric matrix then $g_{t}=e^{t A}$ defines a one parameter curve in $G$, and if $\pi$ is a unitary representation on a complex vector space $V$, then $U_{t}=\pi\left(g_{t}\right)=e^{i t \sigma(A)}$ for some skew symmetric $\sigma(A)$. This way we get a map $A \rightarrow \sigma(A)$ from the space of real skewsymmetric $3 \times 3$ matrices into complex skewhermitian matrices on $V$.

The way to understand this map is to think of $G$ as three dimensional manifold and the vector space of real skewsymmetric $3 \times 3$ matrices as the tangent space at $e$. In fact there are global vector fields acting on functions defined on $G$ corresponding to any skew symmetric $A$,

$$
\left(X_{A}\right) f(g)=\left.\frac{d}{d t} f\left(g e^{t A}\right)\right|_{t=0}
$$

Then

$$
\sigma(A)=\left(X_{A}\right) \pi(e)
$$

and from the representation property

$$
\begin{gathered}
\left(X_{A}\right) \pi(g)=\pi(g) \sigma(A) \\
X_{A} X_{B}=\sigma(A) \sigma(B)
\end{gathered}
$$

The Poisson bracket $\left[X_{A}, X_{B}\right]=X_{A} X_{B}-X_{B} X_{A}$ is to equal $X_{[A B-B A]}$ and we get this a way a representation $\sigma$ of the "Lie Algebra" of $3 \times 3$ skewsymmetric matices in the space of skewhermitian trnasfromations on $V$. Moreover $\sigma([A, B])=[\sigma(A), \sigma(B)] . G$ acts irreducibly on $V$ if and only if $\sigma(A)$ acts irreducibly. We pick a basis $A_{1}, A_{2}, A_{3}$ where

$$
A_{1}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 1 \\
0 & -1 & 0
\end{array}\right] \quad A_{2}=\left[\begin{array}{ccc}
0 & 0 & -1 \\
0 & 0 & 0 \\
1 & 0 & 0
\end{array}\right] \quad A_{3}=\left[\begin{array}{ccc}
0 & 1 & 0 \\
-1 & 0 & 0 \\
0 & 0 & 0
\end{array}\right]
$$

Let us note that

$$
\left[A_{1}, A_{2}\right]=-A_{3},\left[A_{2}, A_{3}\right]=-A_{1},\left[A_{3}, A_{1}\right]=-A_{2}
$$

If we define $\sigma\left(A_{1}\right)=H$ and $Z_{1}=\sigma\left(A_{2}\right)+i \sigma\left(A_{3}\right), Z_{2}=\sigma\left(A_{2}\right)-i \sigma\left(A_{3}\right)$, we can calculate

$$
\begin{aligned}
& {\left[H, Z_{1}\right]=\sigma\left(\left[A_{1}, A_{2}\right]\right)+i \sigma\left(\left[A_{1}, A_{3}\right]\right)=-\sigma\left(A_{3}\right)+i \sigma\left(A_{2}\right)=i Z_{1}} \\
& {\left[H, Z_{2}\right]=\sigma\left(\left[A_{1}, A_{2}\right]\right)-i \sigma\left(\left[A_{1}, A_{3}\right]\right)=-\sigma\left(A_{3}\right)-i \sigma\left(A_{2}\right)=-i Z_{2}}
\end{aligned}
$$

$H$ being skewhermitian on $V$, it has purely imaginary eigenvalues and a complete set of eigenvectors. Let $V=\oplus_{\lambda} V_{i \lambda}$ be the decomposition of $V$ into eigenspaces of $H$. Moreover $e^{2 \pi H}=\pi\left(e^{2 \pi A_{1}}\right)=\pi(e)=I$ The values $\lambda$ are therefore all integers. If $H v=i \lambda v$, then $H Z_{1} v=Z_{1} H v+\left[H, Z_{1}\right] v=$ $i \lambda Z_{1} v+i Z_{1} v=i(\lambda+1) Z_{1} v$. Therefore $Z_{1}$ maps $V_{i \lambda} \rightarrow V_{i(\lambda+1)}$ and similarly $Z_{2}$ maps $V_{i \lambda} \rightarrow V_{i(\lambda-1)}$. It is clear that if we start with some $v_{0} \in V_{i \lambda}$ then $v_{0},\left\{Z_{1}^{k} v_{0}: k \geq 1\right\},\left\{Z_{2}^{k} v_{0}: k \geq 1\right\}$ are all mutually orthogonal. Since the space is finite dimensional, $Z_{1}^{r} v_{0}=Z_{2}^{s} v_{0}=0$ for some $r, s$. If we take $r, s$ to be the smallest such values, then the subspace generated by them has dimension $r+s-1$ and is invariant under $H, Z_{1}, Z_{2}$. Since the representation is irreducible, this must be all of $V$. Another piece of information is that $H$ and $-H$ are conjugate. The set of $\lambda$ 's is therefore symmetric around the
origin. Hence $V$ is odd dimensional and is $\{\lambda\}=\{-k, \ldots, 0, \ldots, k\}$ for some integer $k \geq 0$. This exhausts all possible irreducible representations in the infinitesimal sense and therefore the set of irreducible representations of $G$ cannot be larger. The character of such a representation if it exists is seen to be

$$
\chi_{k}(g)=\hat{\chi}_{k}(\theta)=\sum_{j=-k}^{k} \exp [i j \theta]
$$

where $1, e^{ \pm i \theta}$ are the eigenvalues of $g$. We will try construct them as the natural action of $G$ on the space of homogeneous harmonic polynomials of degree $k$. This dimension is calculated as $\frac{(k+1)(k+2)}{2}-\frac{k(k-1)}{2}=2 k+1 . H$ which is the infinitesimal rotation around $x$-axis is calculated as

$$
H=z \frac{\partial}{\partial y}-y \frac{\partial}{\partial z}
$$

The polynomials $p_{k}^{ \pm}=(y \pm i z)^{k}$ are harmonic in two and therefore three variables and $H p_{k}^{ \pm}= \pm i k p_{k}^{ \pm}$. Therefore this representation has the eigenvlaues $\pm i k$ for $H$ and cannot be decomposed totally in terms of representations of dimension $(2 k-1)$ or less. On the other hand its dimension is only $(2 k+1)$. This is it.

Since we know that $\int_{G} \chi_{k}(g) \chi_{\ell}(g) d g=\delta_{k, \ell}$ it is convenient to determine the weight $w(\theta)$ on $[0, \pi]$ such that it is the probability density of $\theta(g)$ of a random $g$ uniformly distributed over $G$. Then

$$
\int_{0}^{\pi} \hat{\chi}_{k}(\theta) \hat{\chi}_{\ell}(\theta) w(\theta) d \theta=\delta_{k, \ell}
$$

In particular, since $\hat{\chi}_{1}(\theta) \equiv 1$, for $k \geq 2$

$$
\int_{0}^{\pi}\left[\hat{\chi}_{k}(\theta)-\hat{\chi}_{k-1}(\theta)\right] w(\theta) d \theta=0
$$

or

$$
w(\theta)=a+b \cos \theta
$$

Normalization of $\int_{0}^{\pi} w(\theta) d \theta=1$ gives $a=\frac{1}{\pi}$. The orthogonality relation $\int_{0}^{\pi} 1 .(1+2 \cos \theta) w(\theta) d \theta=0$ provides $a+b=0$ or

$$
w(\theta)=\frac{1-\cos \theta}{\pi}
$$

